

# Proton decay matrix elements on lattice

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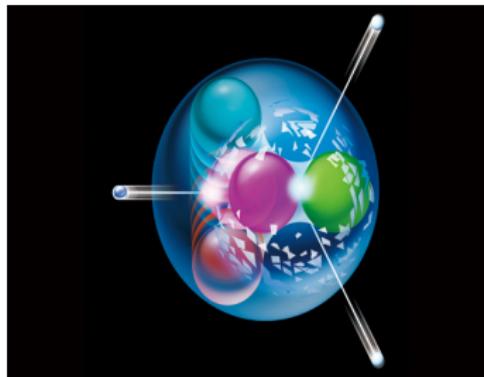
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BNL, Sept. 23-25, 2019



# Introduction

## Proton Decay



**Process**  $p \rightarrow \Pi + \ell$

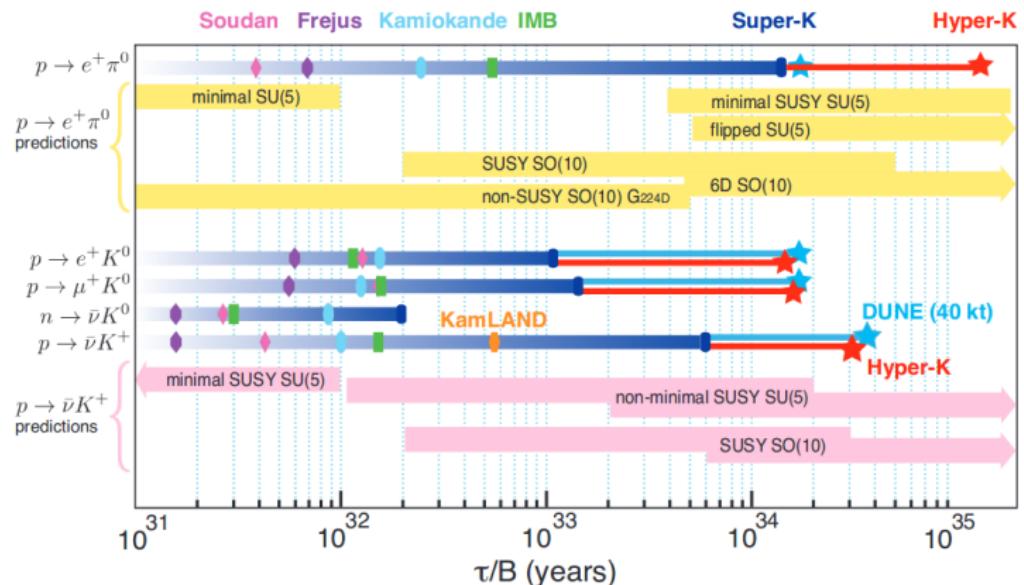
## Baryon Number Violation

**Haven't been observed.**

Lifetime >  $10^{34}$  years

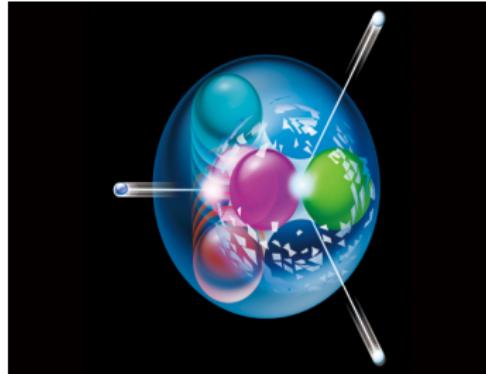
## Does proton even decay?

# Experimental bound



Current proton decay bound in SK, (ABE et al., 2018)

# Proton Decay



**Process  $p \rightarrow \Pi + \ell$**

**Baryon Number Violation**

**Haven't been observed.**

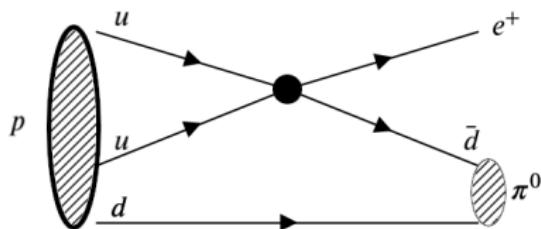
Lifetime  $> 10^{34}$  years

**Motivated by Baryon Asymmetry**

**Possible explanation by**

GUT, SUSY-GUT

# Effective operators



Four-fermion effective operators

$$p \longrightarrow \Pi + \ell$$

Hadronic states : Nonperturbative computation required.

# Effective operator

# GUT, SUSY-GUT

## GUT

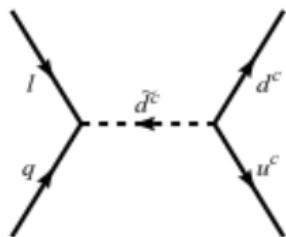
Symmetry group to be  $G \supset SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$

- Coupling unification
- Baryon asymmetry

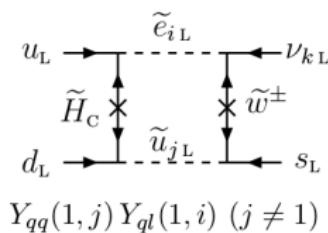
## SUSY-GUT

- Superpartners to particles
- Better unification at higher scale

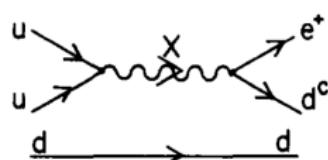
# GUT,SUSY-GUT



(a) d=4 operator



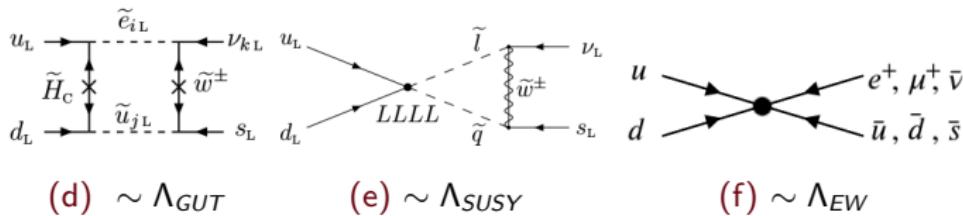
(b) d=5 operator



(c) d=6 operator

Possible BV operators in (SUSY-)GUT

# GUT,SUSY-GUT



Proton decay operator at different scales

Model parameters come into Wilson coefficients

- (a)  $Y_{qq}, Y_{ql}, Y_{ud}, Y_{ue}$
- (b)  $M_{H_C}$
- (c)  $m_{\tilde{l}}, m_{\tilde{q}}$ , triangle loop integrals, ...

# Effective operators

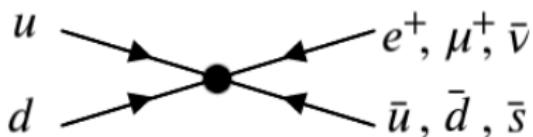


Figure 1: Four-fermion effective operators

Effective operator :  $\mathcal{O}_{\Gamma\Gamma'} = (qq)_\Gamma(q\ell)_{\Gamma'}$ ,

$(XY)_\Gamma = (X^T \mathcal{C} P_\Gamma Y)$        $\mathcal{C}$  := (Charge Conjugation Matrix)

$\langle \Pi \bar{\ell} | p \rangle_{GUT} \sim C^{\Gamma\Gamma} \langle \Pi \bar{\ell} | \mathcal{O}_{\Gamma\Gamma'} | p \rangle_{SM} = C^{\Gamma\Gamma} \bar{v}_\ell \langle \Pi | (qq)_\Gamma P_{\Gamma'} q | p \rangle$ ,

where  $C^{\Gamma\Gamma'}$  is a wilson coefficient,  $\Pi$  is a meson, and  $p$  is a proton.



# Decay rate

The decay rate  $\Gamma$  is calculated from the hadronic matrix element,

$$\begin{aligned} & \langle \Pi(p') | O^{\Gamma\Gamma'}(q) | N(p, s) \rangle \\ &= \bar{v}_\ell P_{\Gamma'} \left[ W_0^{\Gamma\Gamma'}(q^2) - \frac{i\cancel{q}}{m_N} W_1^{\Gamma\Gamma'}(q^2) \right] u_N(p, s) \quad (1) \\ &= \bar{v}_\ell P_{\Gamma'} W_0^{\Gamma\Gamma'}(q^2) u_N(p, s) + O(m_I/m_N) \bar{v}_\ell u_N(p, s) \end{aligned}$$

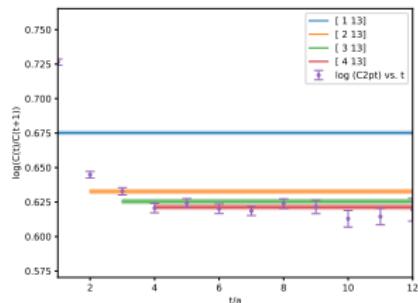
where  $\Pi$  a meson,  $N$  a nucleon, and  $W_{0,1}$  decay form factor(AOKI et al., 2000). Then the decay rate is

$$\Gamma(p \rightarrow \Pi + \bar{\ell}) = \frac{(m_p^2 - m_\Pi^2)^2}{32\pi m_p^3} \left| \sum_I C_I W_0^I(p \rightarrow \Pi + \bar{\ell}) \right|^2. \quad (2)$$

# Matrix elements on lattice

## Hadronic states

## Correlation function



$$\begin{aligned}
 & C_K^{2pt}(t, \vec{p}) \\
 &= \sum_{\vec{x}} e^{i\vec{p} \cdot \vec{x}} \langle 0 | J_K(t, \vec{x}) J_K^\dagger(0, \vec{0}) | 0 \rangle \\
 &= \text{asymptotic states} + ...
 \end{aligned}$$

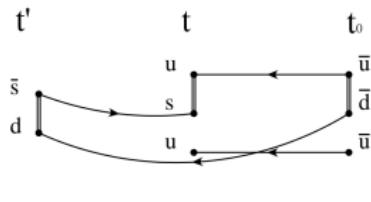
$$p_K = [0, 1, 1] p_{\min}$$

## Excited states at early times

## Ground state at late times

# Three-point function

(Meson)-(Decay Operator)-(Proton)



$$\begin{aligned} C^{3pt}(t, t') &= \sum_{\vec{x}, \vec{x}'} e^{i(\vec{p}' \cdot \vec{x}' - \vec{q} \cdot \vec{x})} \langle 0 | J_\Pi(x') \mathcal{O}(x) \bar{J}_N(x_0) | 0 \rangle \\ &= \langle \Pi(\vec{p}') | \mathcal{O} | N(\vec{p}) \rangle \\ &\times \frac{C_\Pi^{2pt}(t' - t, \vec{p}')}{\sqrt{Z_\Pi}} \frac{\text{Tr}[P C_p^{2pt}(t, \vec{p})]}{\sqrt{Z_p}} \end{aligned}$$

# Lattice settings

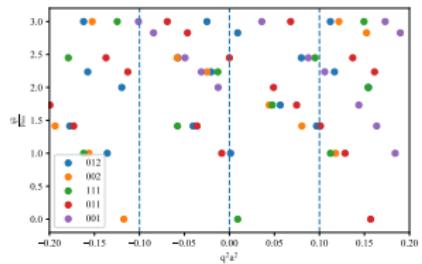
## Lattice parameters

lattice size	$24^3 \times 64 \times 24$
gauge action	Iwasaki-DSDR
fermion	DWF
$\beta$	1.633
lattice cutoff	$a^{-1} = 1\text{GeV}$
$m_f a$	0.00107
$m_h a$	0.0850
$m_\pi a$	0.1387
$m_K a$	0.5051
$m_{\text{res}}$	0.00228
$m_\pi L$	3.3
Deflated CG	2000+1000
AMA	32+1
$N_{cfg}$	102

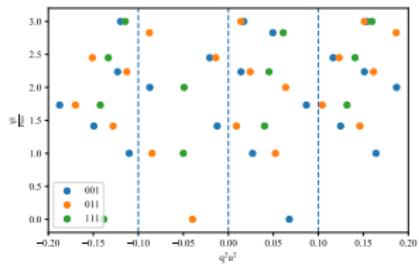
lattice size	$32^3 \times 64 \times 32$
gauge action	Iwasaki-DSDR
fermion	DWF
$\beta$	1.75
lattice cutoff	$a^{-1} = 1.37\text{GeV}$
$m_f a$	0.0001
$m_h a$	0.0450
$m_\pi a$	0.1046
$m_K a$	0.3602
$m_{\text{res}}$	
$m_\pi L$	3.3
Deflated CG	2000+250
AMA	32+1
$N_{cfg}$	20(and counting)

# Kinematic Choice

Energy-momentum conservation &  $q^2 \sim 0$

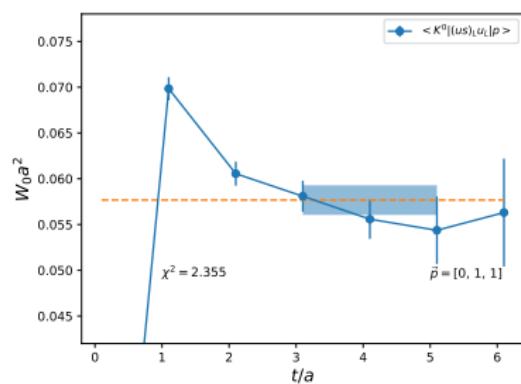


$\pi$



K

# Matrix elements



## decay form factor

$W_0^{LL}(p \rightarrow K^0 e^+)$  at  $t_{sep} = 8$

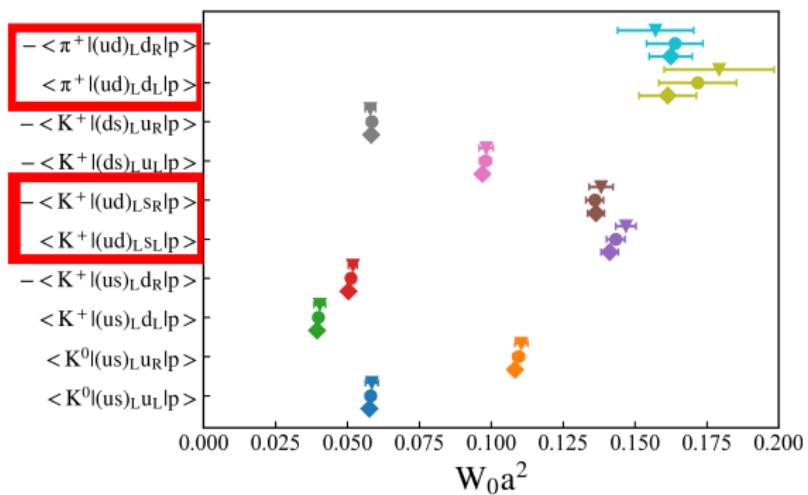
$$p_p = [0, 0, 0] p_{\min},$$

$$p_K = [0, 1, 1] p_{\min}$$

plateau fit  $t=3-5$

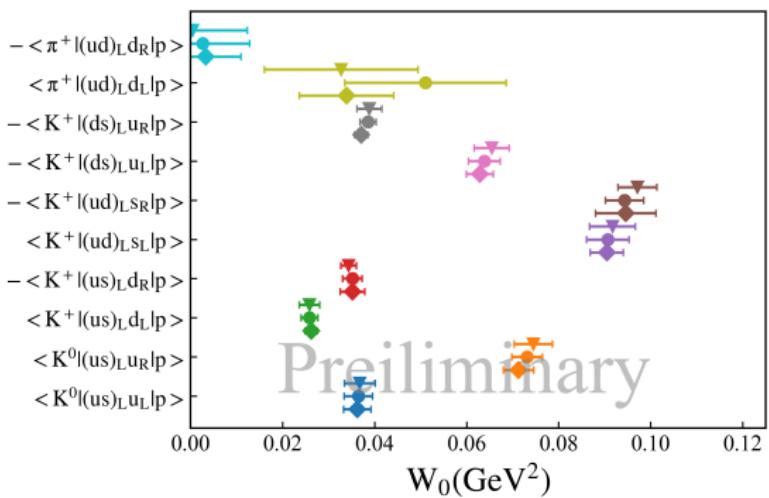
AMA 32+1, 102 configs

# Matrix elements



Decay matrix elements w/ different src-sink separation {8,9,10}

# Matrix elements



Decay matrix elements w/ different src-sink separation {8,9,10}

# Comparison with earlier work

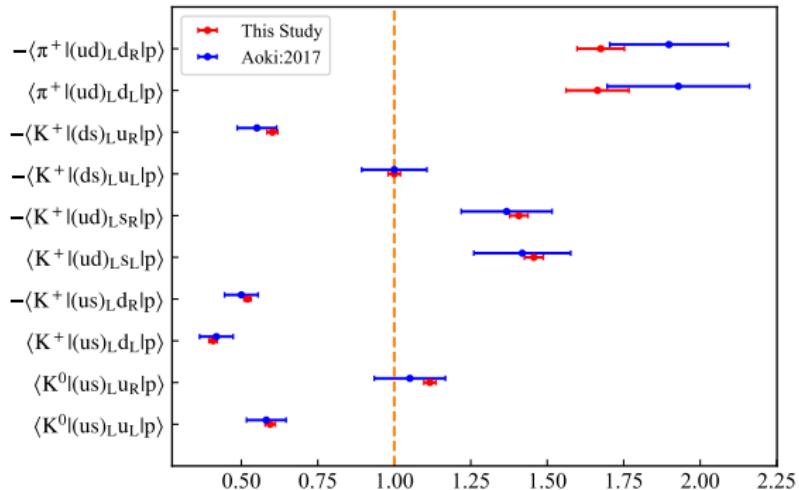
Bare value, but multiplicative renormalization only

→ ratio can be compared with renormalized values

$$W_0^{norm} = \left| \frac{W_0^{\Gamma\Gamma'}(\text{Channel})}{W_0^{\Gamma\Gamma'}(\langle K^+ | (ds)_\Gamma u_{\Gamma'} | p \rangle)} \right| \quad (3)$$

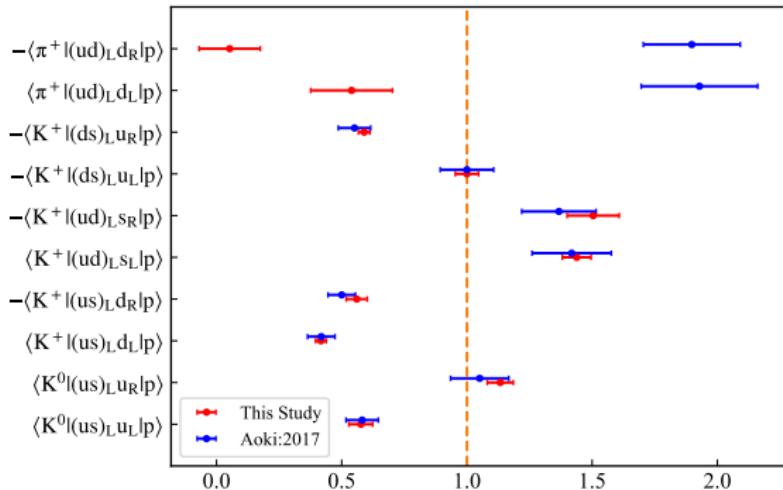


# Comparison with earlier work



Comparison with earlier study, (AOKI et al., 2017)

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Comparison with earlier study, (AOKI et al., 2017)

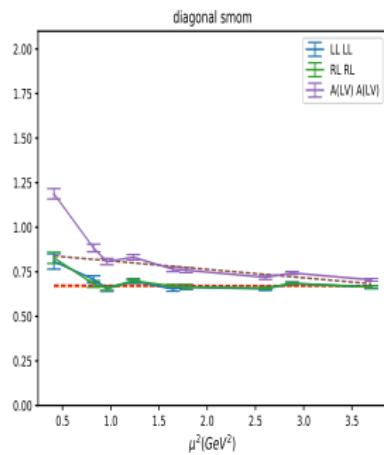
# Comparison to earlier work

	Stat. [%] (This study)	Stat. [%] (Aoki:2017)	Chiral extrapol. [%]	$a^2$ [%]	$\Delta_Z$ [%]
$\langle K^0   (us)_L u_L   p \rangle$	2.80	3.5	3.1	5.0	8.1
$\langle K^0   (us)_L u_R   p \rangle$	1.77	2.8	2.8	5.0	8.1
$\langle K^+   (us)_L d_L   p \rangle$	3.32	4.4	7.5	5.0	8.1
$-\langle K^+   (us)_L d_R   p \rangle$	2.24	3.7	3.5	5.0	8.1
$\langle K^+   (ud)_L s_L   p \rangle$	2.13	3.0	3.9	5.0	8.1
$-\langle K^+   (ud)_L s_R   p \rangle$	2.12	3.2	1.6	5.0	8.1
$-\langle K^+   (ds)_L u_L   p \rangle$	2.01	2.8	2.1	5.0	8.1
$-\langle K^+   (ds)_L u_R   p \rangle$	2.96	3.6	2.7	5.0	8.1
$-\langle \pi^+   (ud)_L d_R   p \rangle$	6.17	3.4	2.7	5.0	8.1
$\langle \pi^+   (ud)_L d_R   p \rangle$	4.62	3.0	2.7	5.0	8.1

Left : Comparison of statistical errors. Right: Systematic errors in chiral extrapolation,  $O(a^2)$ ,  $\Delta_Z$  ( (AOKI et al., 2017))



# Renormalization



Three quark op. renormalization

32+1 AMA, 20 configs, Landau gauge

RI-SMOM scheme

Two-loop Matching (GRACEY, 2012)

fit region  $(pa)^2 = 1 - 3.5 \text{ GeV}^2$

$$U_{LL}^{\overline{MS} \leftarrow \text{latt}}(\mu = 2 \text{ GeV}) = 0.67(2)$$

$$U_{RL}^{\overline{MS} \leftarrow \text{latt}}(\mu = 2 \text{ GeV}) = 0.68(1)$$

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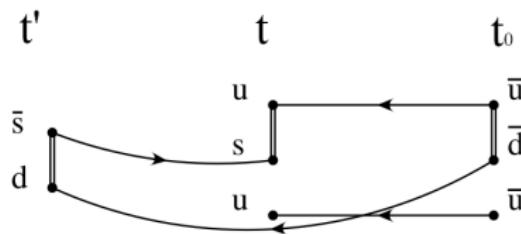
## Future projects

# Future projects

Proton decay matrix elements can be investigated further to see:

- Induced Nucleon Decay from Dark matter
- Vector meson channels from proton decay
- Three body decay channel from proton decay

# Vector meson channels



- Same computation with different  $\Gamma$  structures
- Different form factor decomposition
- Asymptotic vector meson channel should be there.

$$\langle K^{*i}(Q)\ell(p')|O_{d=6}|p(p,s)\rangle = \epsilon_\mu^i \bar{v}_\ell^c [F_1 \gamma_5 \gamma^\mu + F_2 i \gamma_5 \sigma^{\mu\nu} Q_\nu + F_3 \gamma_5 Q^\mu$$

$$+ F'_1 \gamma^\mu + F'_2 i \sigma^{\mu\nu} Q_\nu + F'_3 Q^\mu] u_N$$



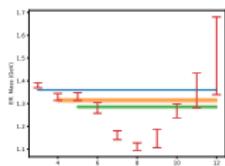
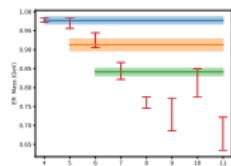
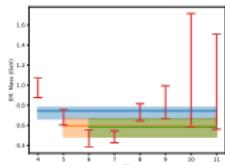
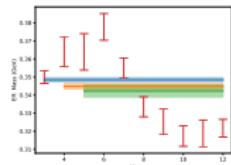
# Vector meson channels

Lattice parameters(BOYLE et al., 2016)

lattice size	$24^3 \times 64 \times 16$
gauge action	Iwasaki
fermion	DWF
$\beta$	2.13
lattice cutoff	$a^{-1} = 1.78\text{GeV}$
$m_I a$	0.0005
$m_h a$	0.04
$m_\pi$	339.6 MeV
$m_\pi L$	4.568



# Vector meson channels



Two point fcn.

$m_\pi = 345(6)\text{MeV}$ , p-val = 0.08

$m_\rho = 916(46)\text{MeV}$ , p-val = 0.00

$m_\sigma = 595(40)\text{MeV}$ , p-val = 0.99

$m_p = 1.315(16)\text{GeV}$ , p-val=0.924

\*No disconnected diagram is included.

# Three body decay

A few reasons to compute three body decay :

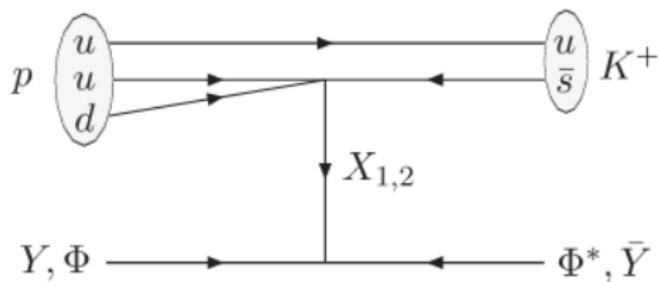
- Resonant to vector meson channels :  $p \rightarrow K^* + \bar{\ell} \rightarrow (K\pi)\bar{\ell}$
- Decay rate ratio  $(\Gamma(p \rightarrow \pi\pi e^+)/\Gamma(p \rightarrow \pi e^+))$  estimates to  $\sim 24\text{--}150\%$  (WISE; BLANKENBECLER; ABBOTT, 1981)
- Prime channel of next generation experiment
- Numerically cheapest among three body decay channels

Decay Mode	Water Cherenkov Efficiency	Cherenkov Background	Liquid Argon TPC Efficiency	Liquid Argon TPC Background
$p \rightarrow K^+\bar{\nu}$	19%	4	97%	1
$p \rightarrow K^0\mu^+$	10%	8	47%	< 2
$p \rightarrow K^+\mu^-\pi^+$			97%	1
$n \rightarrow K^+e^-$	10%	3	96%	< 2
$n \rightarrow e^+\pi^-$	19%	2	44%	0.8

DUNE proton decay efficiency

# Future projects

Induced Nucleon Decay model (DAVOUDIASL et al., 2010)



- DM can annihilate the nucleon
- $\mathcal{L}_{eff} \sim (1/\Lambda^3) u_R d_R d_R Y_R \Phi + h.c.$
- $\langle \Pi(p') | O^{\Gamma\Gamma'}(q) | N(p, s) \rangle = P_{\Gamma'} \left[ W_0^{\Gamma\Gamma'}(q^2) + \frac{m_Y}{m_N} W_1^{\Gamma\Gamma'}(q^2) \right] u_N(p, s)$

# Conclusion

## Conclusion

- Proton decay matrix elements on the two lattice ensemble with chiral fermions at physical scale
  - Three quark op. non perturbative renormalization
  - New channels(vector meson, three body)
  - Induced nucleon decay by DM

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## Reference

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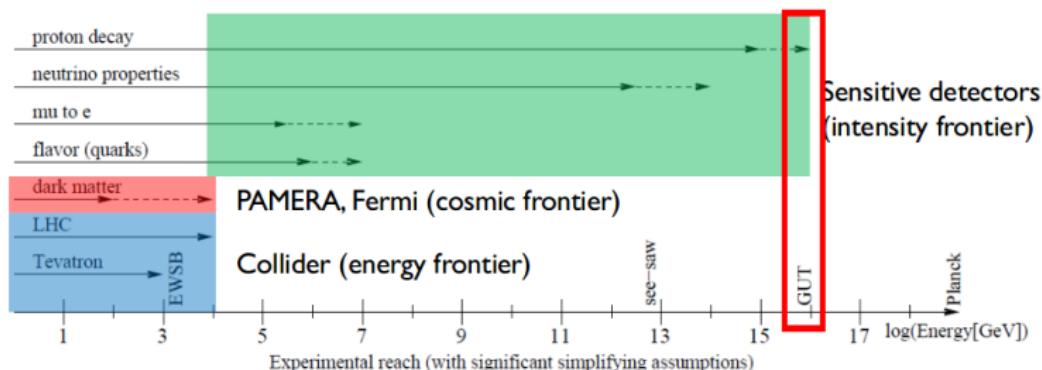


Figure 2: Energy scale of search, Zoltan Ligeti

## Baryon asymmetry

Nonzero net baryon number

$$\frac{n_B - \bar{n}_B}{n_\gamma} \sim 10^{-10}$$

### Sakharov's conditions

- At least one B violating process
- C- and CP-violation
- interactions outside of thermal equilibrium

# Matrix elements

Define the ratio  $R_3(t, t') = \frac{C^{3pt}(t, t')}{C_{\Pi}^{2pt}(t' - t, \vec{p}') \text{Tr}[P C_p^{2pt}(t, \vec{p})]} \sqrt{Z_{\Pi}} \sqrt{Z_p}$ .

As  $t \rightarrow \infty$ ,  $R_3(t, t') \rightarrow \langle \Pi(p') | O^{\Gamma\Gamma'}(q) | N(p, s) \rangle$ , giving decay form factors  $W_{0,1}(q^2)$

$$\text{Tr}[R_3 P_L P_4] = W_0^{\Gamma L}(q^2) - \frac{i q_4}{m_N} W_1^{\Gamma L}(q^2).$$

$$\text{Tr}[R_3 P_L i P_4 \gamma_j] = \frac{q_j}{m_N} W_1^{\Gamma L}(q^2)$$

Momentum transfer is chosen to be  $q^2 \sim 0$  :  $\vec{p} = \vec{q} + \vec{p}'$

# Other computations

		$ \alpha  \text{ [GeV}^3]$	$ \beta  \text{ [GeV}^3]$	
	Donoghue and Goldwicht [4]	0.003		Bag model
	Thomas and McKellar [7]	0.02		Bag model
	Meljanac et al. [5]	0.004		Bag model
QCD model	Ioffe [2]	0.009		Sum rule
calculation	Krasnikov et al. [6]	0.003		Sum rule
	Ioffe and Smilga [8]	0.006		Sum rule
	Tomozawa [3]	0.006		Quark model
	Brodsky et al. [9]	0.03		
	Hara et al. [10]	0.03		WF, $a = 0.11$ fm
	Bowler et al. [11]	0.013	0.010	WF, $a = 0.22$ fm
Lattice QCD	Gavela et al. [12]	0.0056(8)	$\simeq  \alpha $	WF, $a = 0.09$ fm
$N_f = 0$	JLQCD [13]	0.015(1)	0.014(1)	WF, $a = 0.09$ fm
	CP-PACS & JLQCD [14]	0.0090(09)( $^{+5}_{-19}$ )	0.0096(09)( $^{+6}_{-20}$ )	WF, continuum limit
	Aoki et al. [15]	0.0100(19)	0.0108(21)	DWF, $a = 0.15$ fm
Lattice QCD	Aoki et al. [15]	0.0118(21)	0.0118(21)	DWF, $a = 0.12$ fm
$N_f = 2$				
Lattice QCD	This work	0.0112(25)	0.0120(26)	DWF, $a = 0.12$ fm
$N_f = 2 + 1$				

Figure 3